

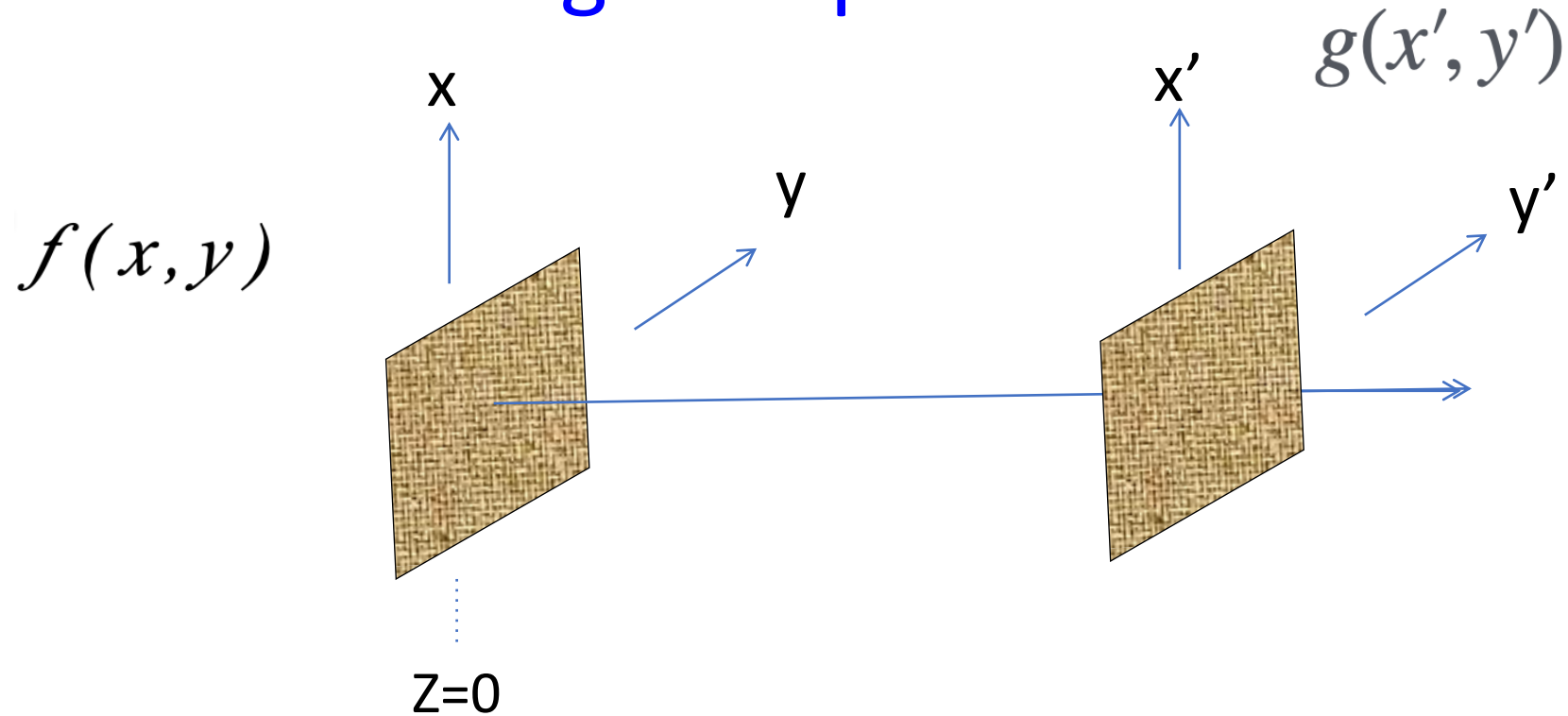
Computational Optical Imaging

Lecture 2

Outline

- Last time: Angular Spectrum Method (ASM) and 2D FTs
- Fresnel diffraction
- Near field and NSOM (near field optical microscopy)
- Far field
- Exercise #1
- Gratings
- Talbot imaging
- Bessel beam and Light sheet imaging
- Exercise #2

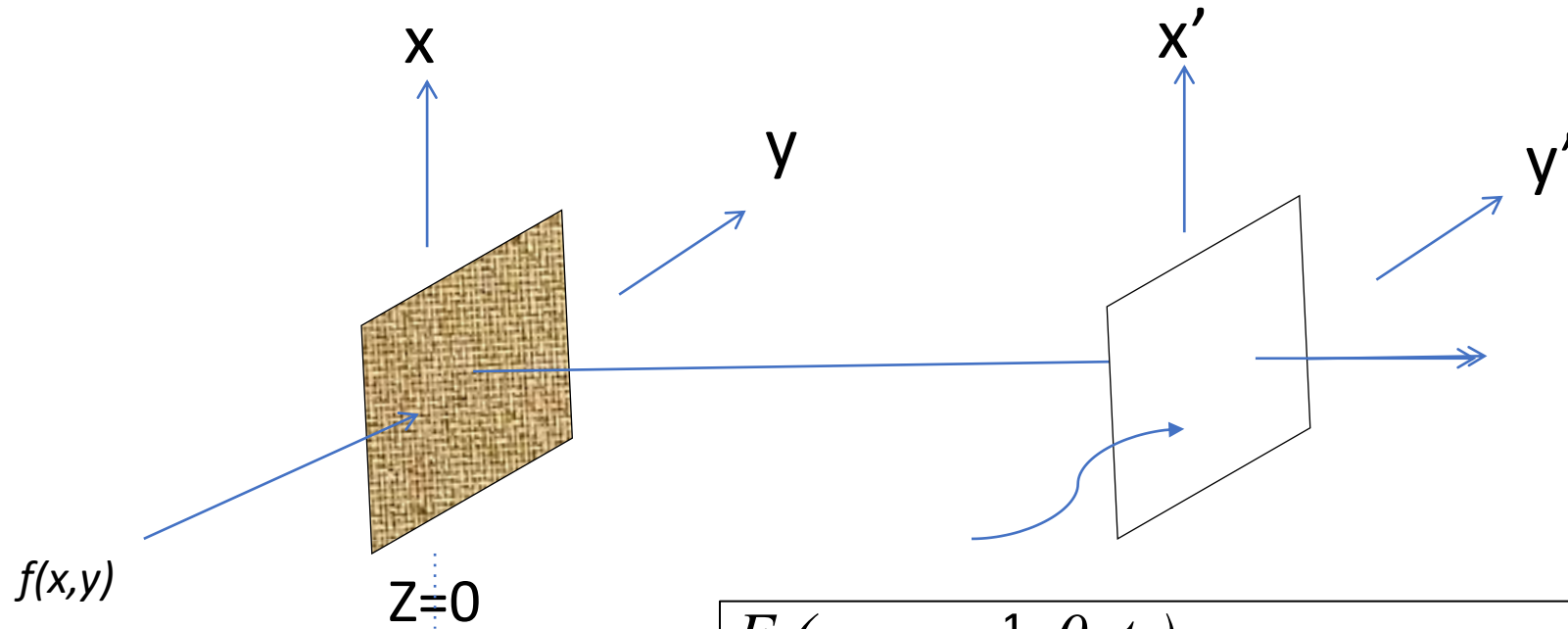
Angular spectrum



3 steps

1. Take the 2D Fourier transform of $f(x, y)$
2. Multiply each spatial frequency component by its eigenvalue
3. Take the inverse 2D Fourier transform to get $g(x', y')$

Angular spectrum



$$E(x, y, z=0, t) = A f(x, y) e^{j\omega t}$$

$$f(x, y) \xleftrightarrow{\text{2D FT}} F(u, v)$$

$$E(x, y, z \neq 0, t) = e^{j\omega t} \iint F(u, v) e^{-j(2pux + 2pvy + \sqrt{k^2 - (2pu)^2 - (2pv)^2} z)} du dv$$

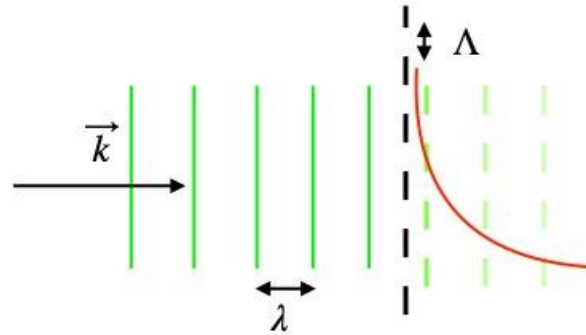
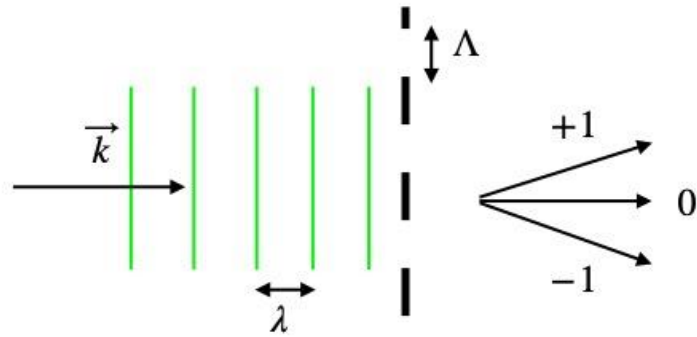
$$E(x, y, z = 0, t) = A e^{j\omega t} f(x, y) = A e^{j\omega t} (2p)^2 \iint F(u, v) e^{-j2p(ux + vy)} du dv$$

$$k_x = 2pu \quad k_y = 2pv \quad k_z = \sqrt{k^2 - k_x^2 - k_y^2} = \sqrt{k^2 - (2pu)^2 - (2pv)^2}$$

Evanescent Waves

$$k_x^2 + k_y^2 + k_z^2 = \omega^2 \mu \epsilon = \left(\frac{2\pi}{\lambda}\right)^2 = k^2$$

$$k_z = \sqrt{k^2 - (k_x^2 + k_y^2)}$$



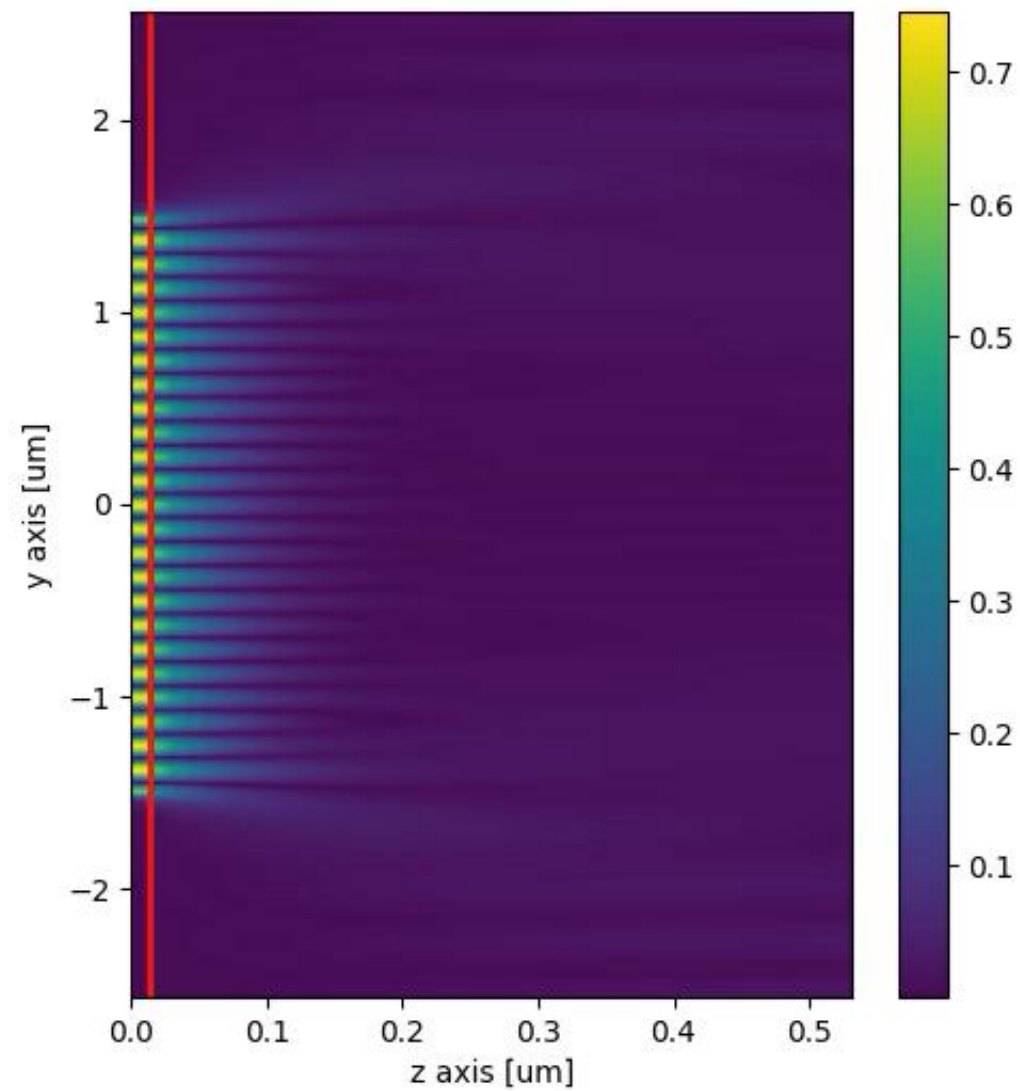
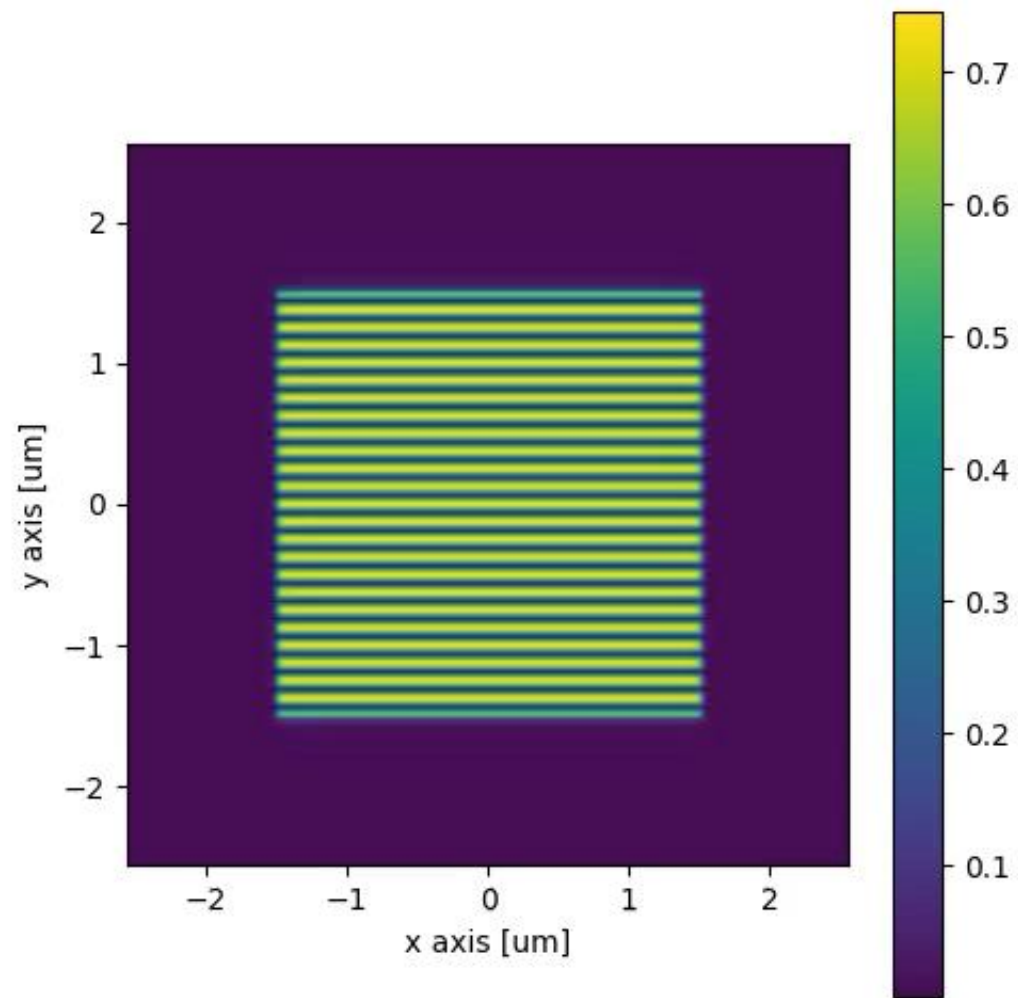
Exponential Decay

$\Lambda < \lambda$

$$k_x = \frac{2\pi \cos \theta_x}{\lambda} = \frac{2\pi}{\Lambda}$$

$$u = \frac{1}{\Lambda}$$

$$k_z = \sqrt{k^2 - (k_x)^2} = \sqrt{\left(\frac{2\pi}{\lambda}\right)^2 - \left(\frac{2\pi}{L}\right)^2}$$



$\lambda=532\text{nm}$. $\Lambda=250\text{nm}$

Near-field imaging NSOM

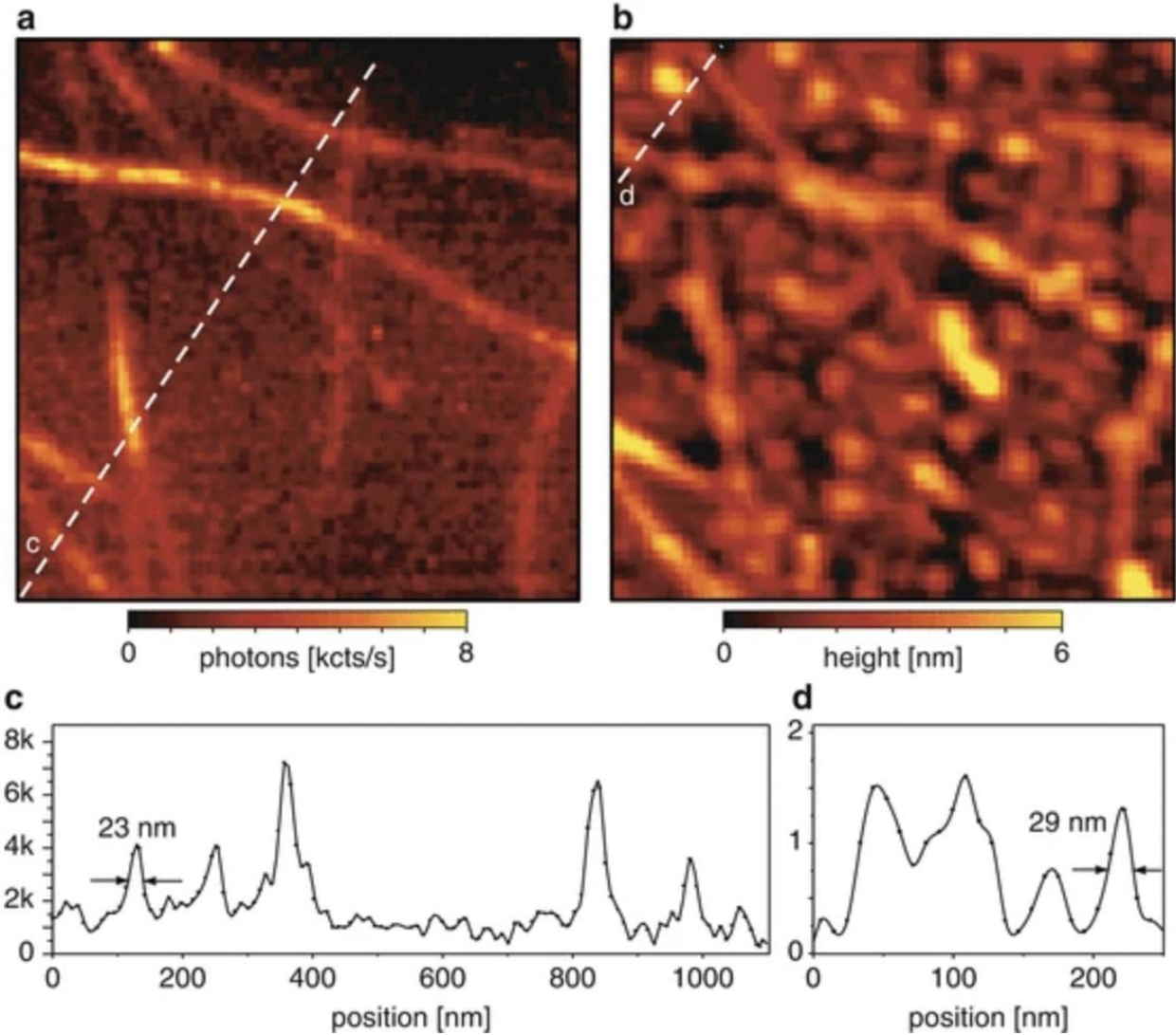
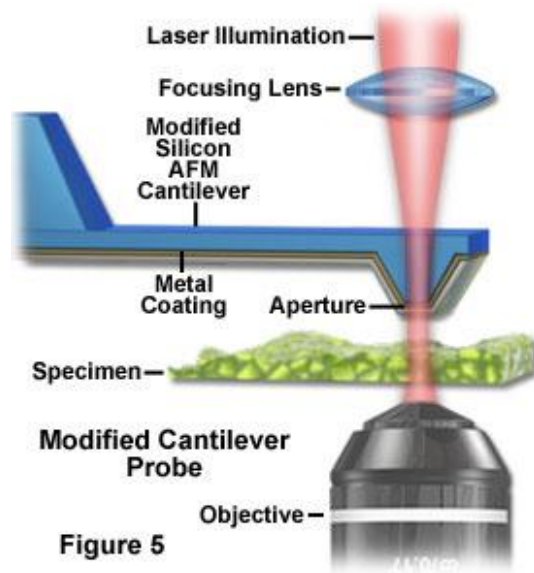
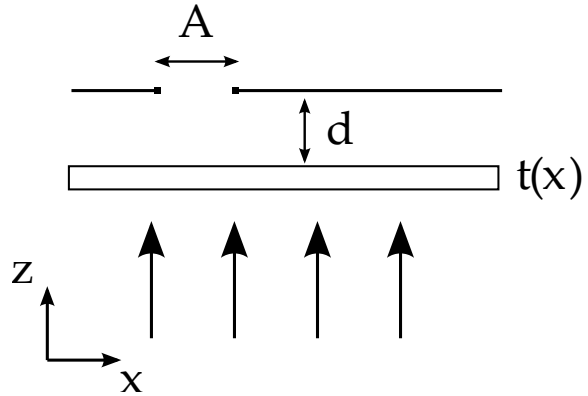
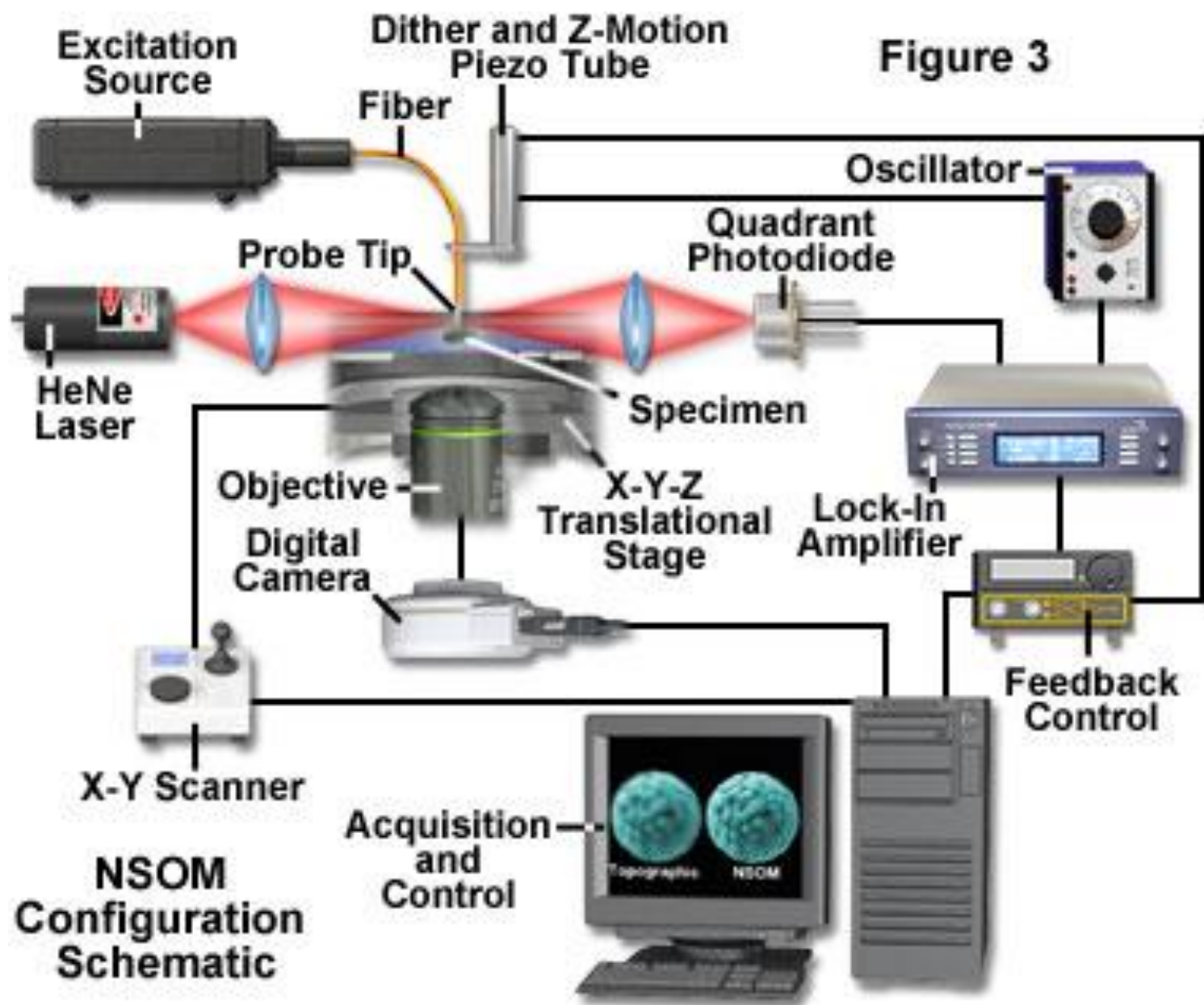
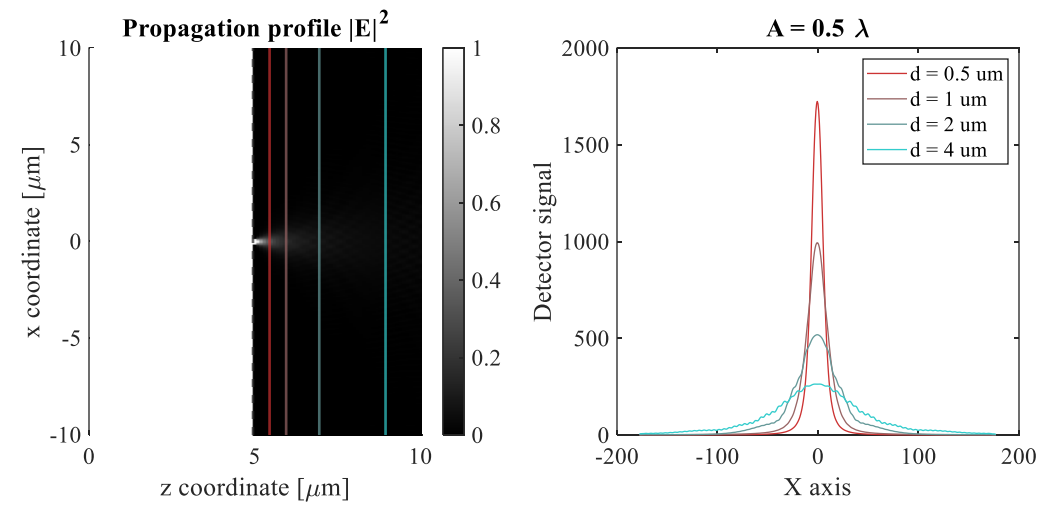
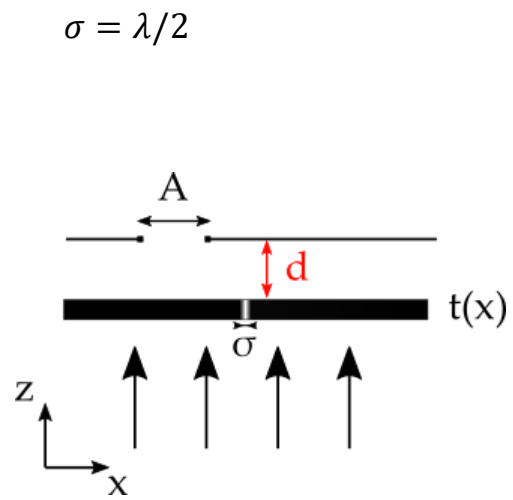


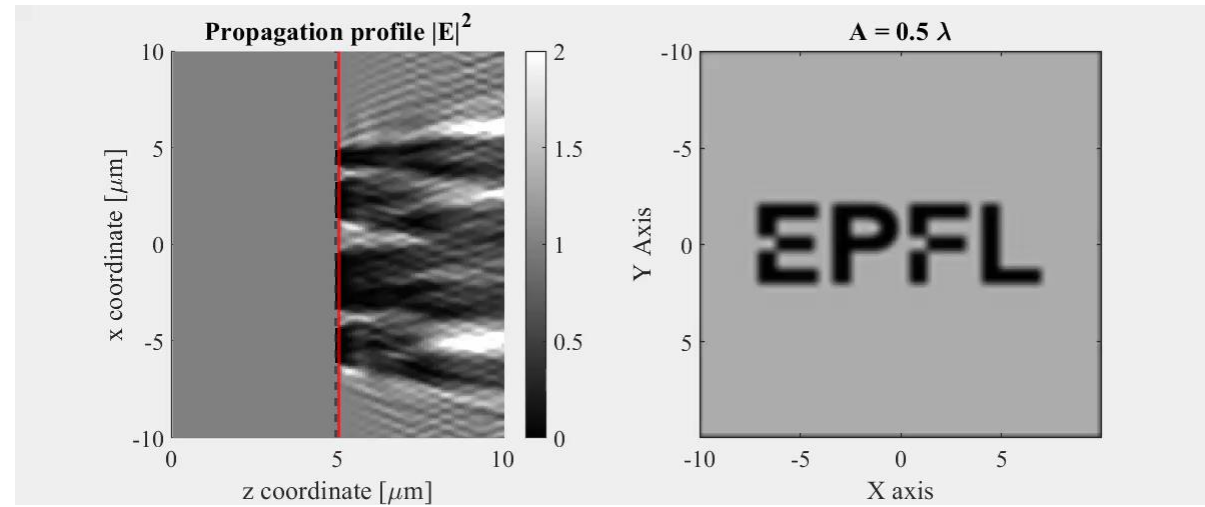
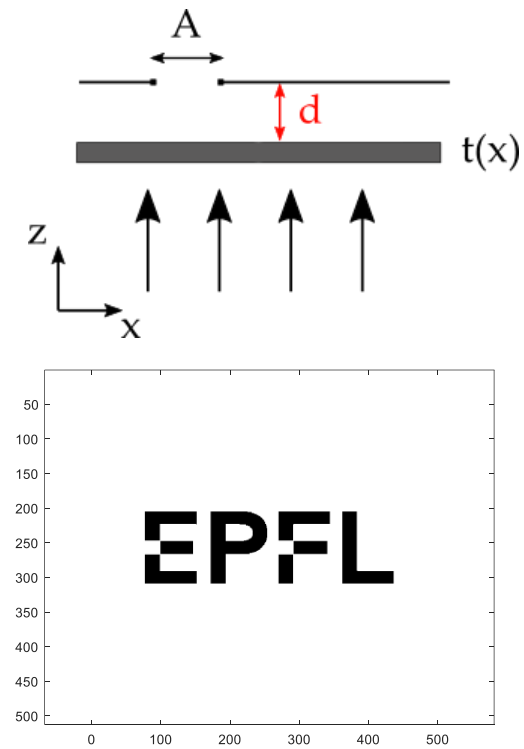
Figure 3



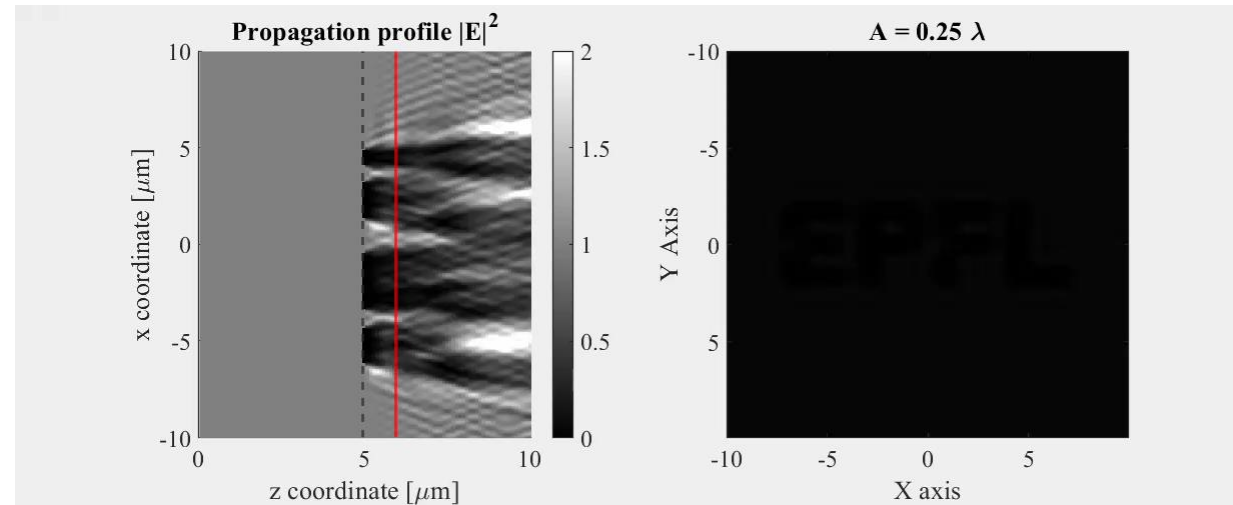
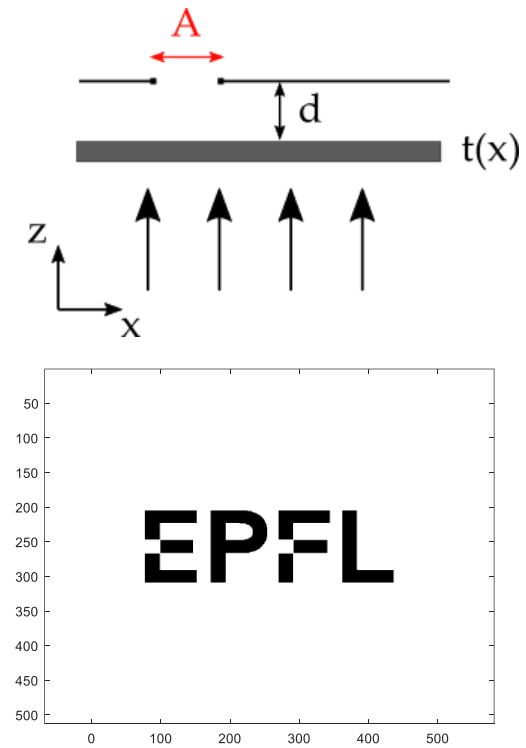
Example 2: Point source



Example 4: Image

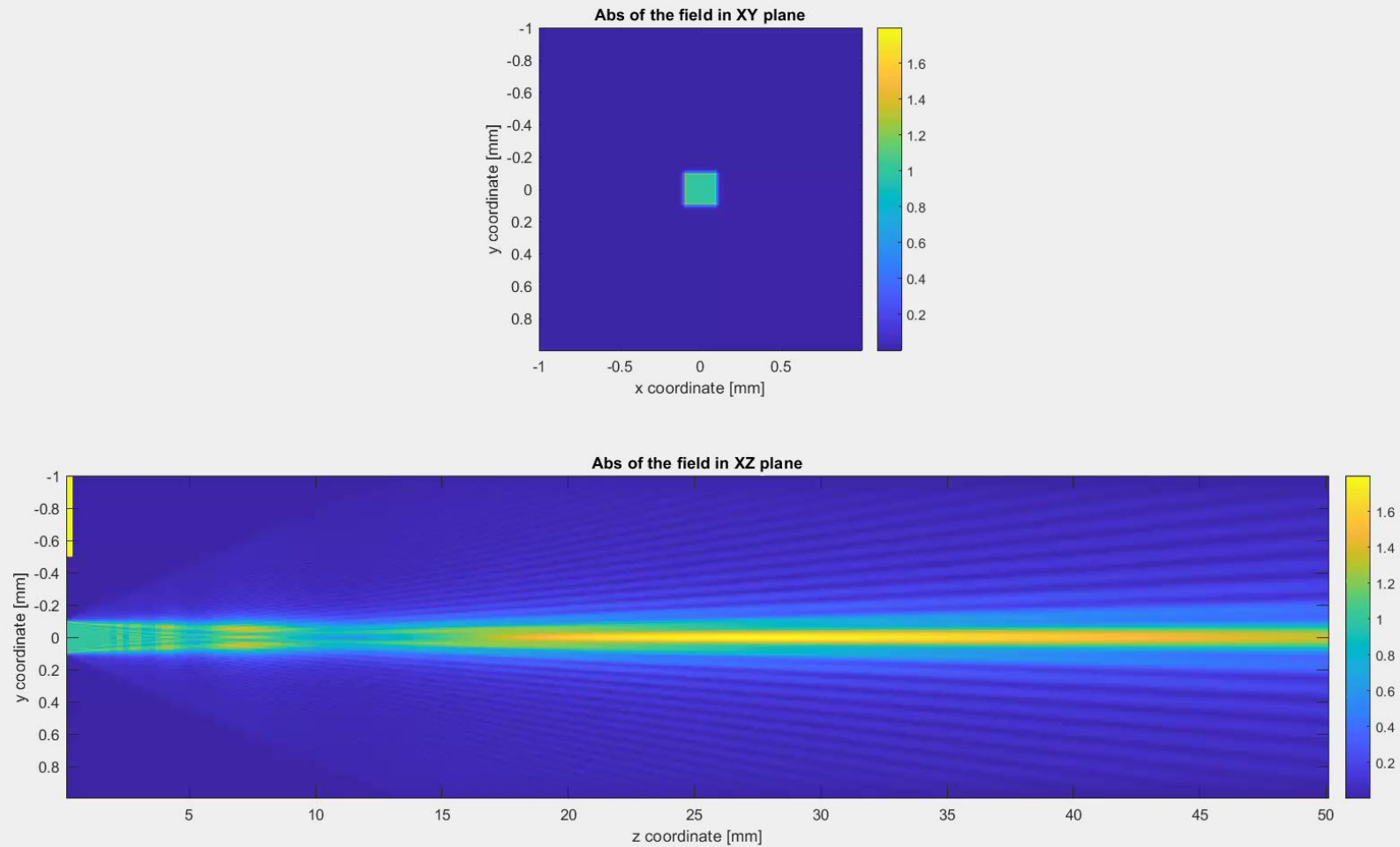


Example 4: Image



Fraunhofer Diffraction or Far field

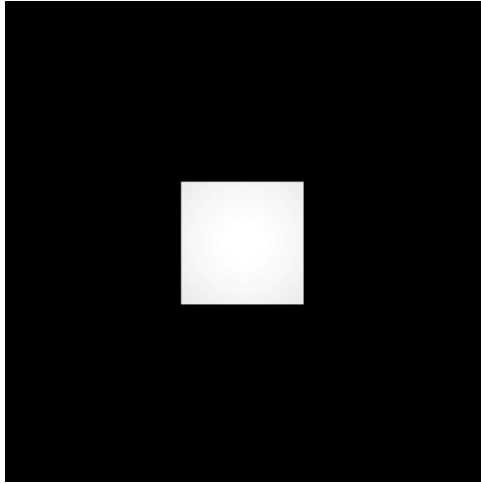
Propagation of square aperture



Square aperture - linear propagation

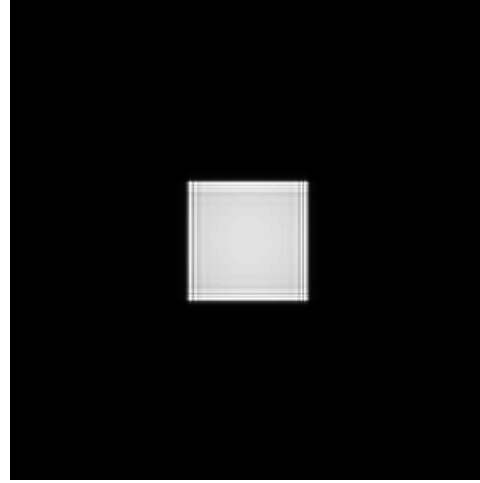
Aperture

$z = 0\text{mm}$

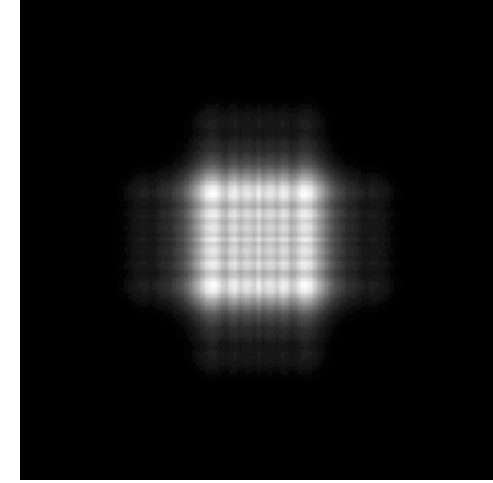


near field

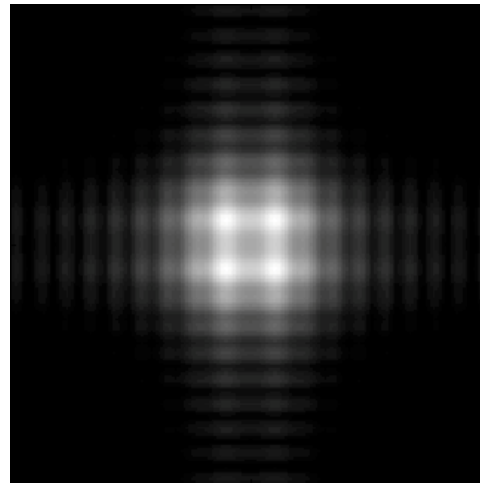
$z = 1\text{mm}$



$z = 20\text{mm}$



$z = 100\text{mm}$



far field (rescaled)

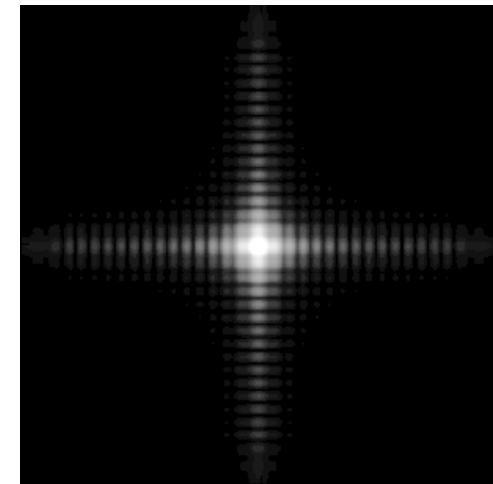
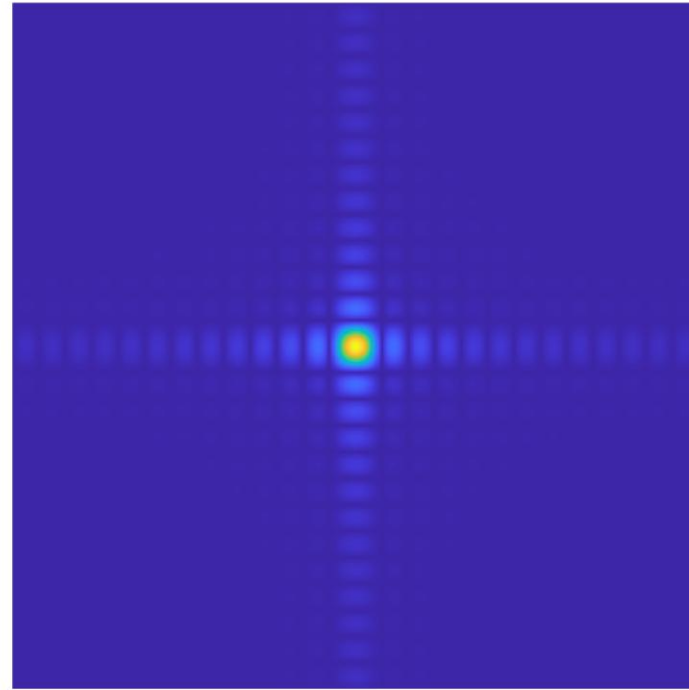
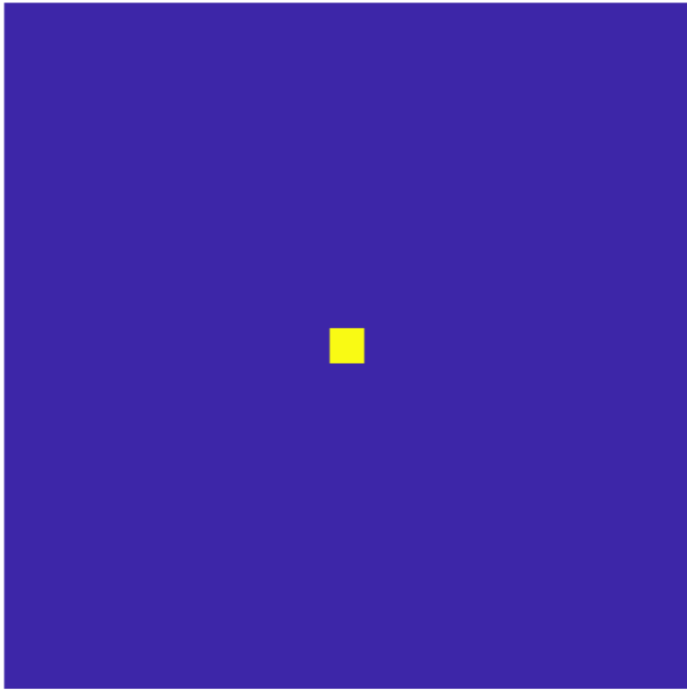


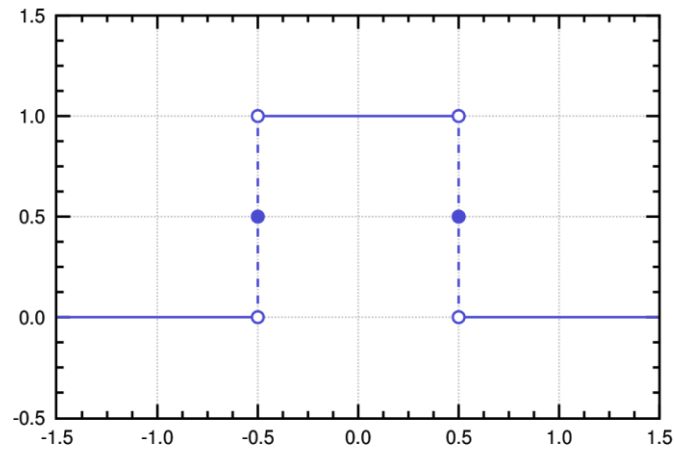
Image size = 2x2mm
Aperture side = 550 μm
Wavelength = 800nm

2D square and 2D sinc

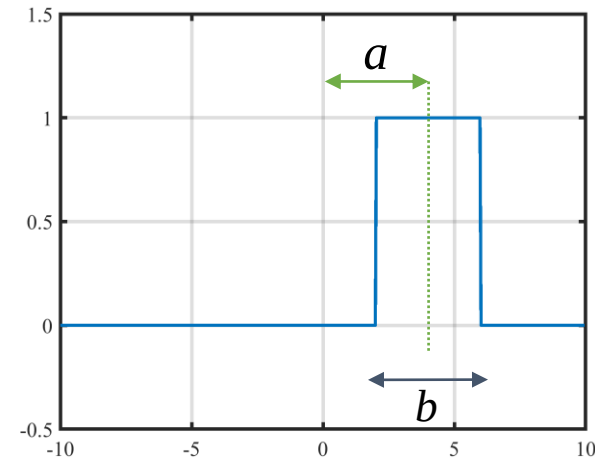


1D rect and 1D sinc

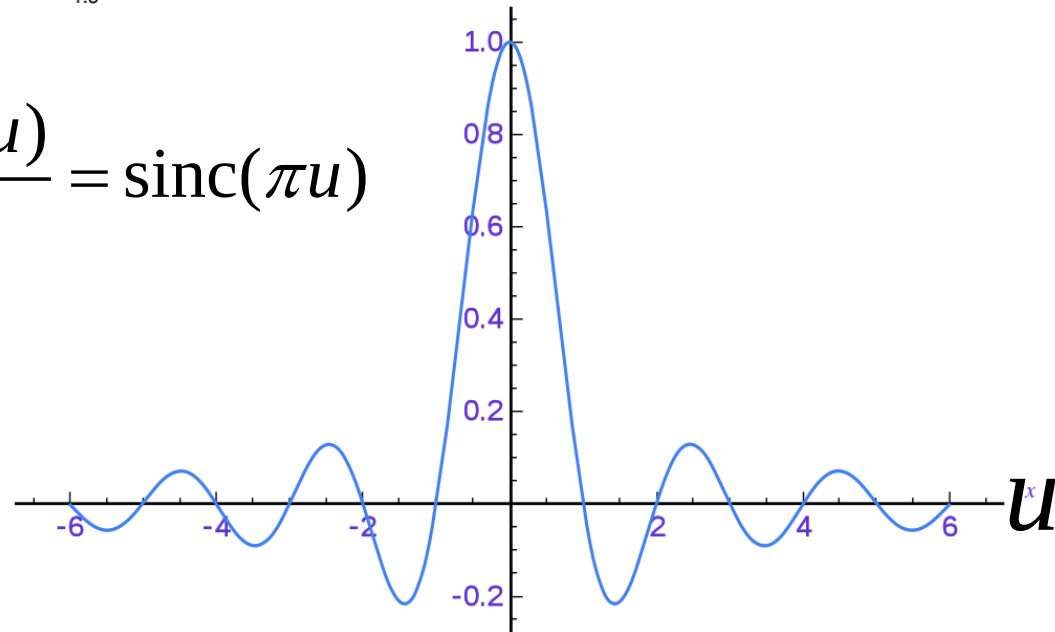
$$\text{rect}(x) = \Pi(x) = \begin{cases} 0, & \text{if } |x| > 1/2 \\ 1/2, & \text{if } |x| = 1/2 \\ 1, & \text{if } |x| < 1/2 \end{cases}$$



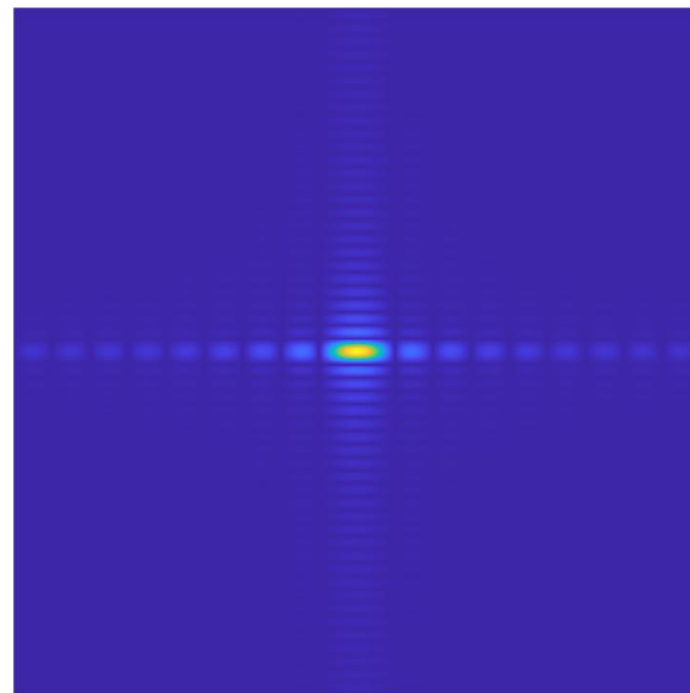
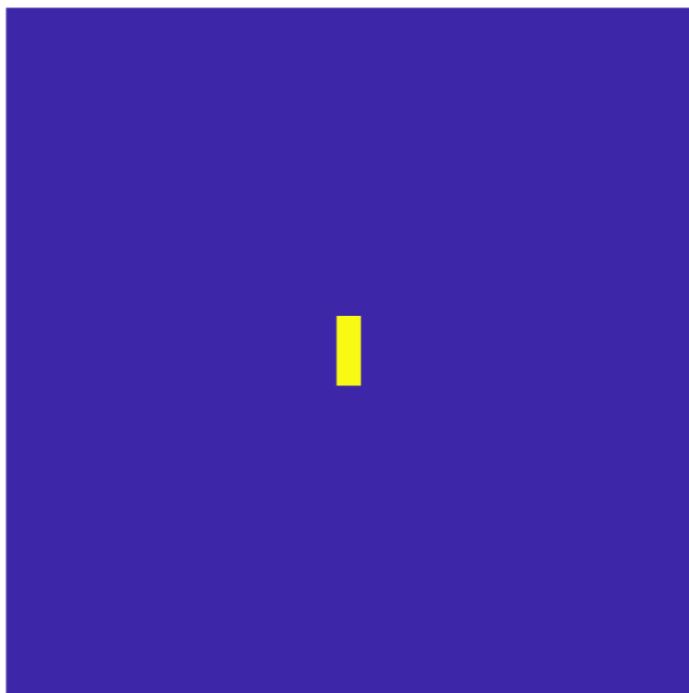
$$\text{rect}[(x-a)/b]$$



$$\int_{-\infty}^{\infty} \text{rect}(x) e^{-j2\pi ux} dx = \frac{\sin(\pi u)}{\pi u} = \text{sinc}(\pi u)$$



2D rectangle



Paraxial Approximation

$$e^{-j(k_x x + k_y y + k_z z)}$$

$$k_x^2 + k_y^2 + k_z^2 = k^2 = \left(\frac{2\pi}{\lambda}\right)^2$$

$$e^{-j(k_x x + k_y y + k_z z)} = e^{-j(k_x x + k_y y)} e^{-j\sqrt{k^2 - k_x^2 - k_y^2} z}$$

$$\cong e^{-j(k_x x + k_y y)} e^{-j\left(k - \frac{k_x^2 + k_y^2}{2k}\right) z} \quad \text{Note: } \sqrt{1-x} \approx 1 - \frac{x}{2}, \quad x \ll 1$$

Paraxial Angular Spectrum

$$E(x, y, z) = \iint F(u, v) e^{-j(2pu x + 2pv y + \sqrt{k^2 - (2pu)^2 - (2pv)^2} z)} du dv$$

Paraxial plane wave:

$$\tilde{A}(k_x, k_y, z) = e^{j \left[\frac{k_x^2 + k_y^2}{2k} \right] z} \tilde{A}(k_x, k_y, z=0)$$

Convolution Theorem: $g(x', y') = \iint f(x, y) h(x' - x, y' - y) dx dy$

$$G(u, v) = F(u, v) H(u, v)$$

G, F and H are Fourier transforms of g, f and h respectively

Note that Fourier Transform of the exponential is:

$$e^{-p(x^2 + y^2)} \Leftrightarrow e^{-p(u^2 + v^2)} \quad \text{where} \quad k_x = 2pu \quad k_y = 2pv$$

Fresnel Diffraction

Let us apply the Convolution Theorem on Beam Propagation Equation:

$$A(x', y', z) = \frac{1}{j\lambda z} \iint A(x, y, z=0) e^{-\frac{j\pi[(x-x')^2 + (y-y')^2]}{\lambda z}} dx dy$$

Recall: $E_x(x, y, z, t) = A(x, y, z) e^{j\omega t} e^{-jkz}$

Finally:

$$E_x(x', y', z, t) = A(x', y', z) e^{j\omega t} e^{-jkz} = \frac{1}{j\lambda z} e^{j\omega t} e^{-jkz} \iint A(x, y, z=0) e^{-\frac{j\pi[(x-x')^2 + (y-y')^2]}{\lambda z}} dx dy$$

Far Field

$$\begin{aligned}
 E_x(x', y', z, t) &= A(x', y', z) e^{j\omega t} e^{-jkz} \\
 &= e^{j\omega t} e^{-jkz} \iint A(x, y, z=0) e^{-\frac{j\rho[(x-x')^2 + (y-y')^2]}{lz}} dx dy \\
 &= e^{j\omega t} e^{-jkz} e^{-\frac{j\rho(x'^2 + y'^2)}{lz}} \iint A(x, y, z=0) e^{-\frac{j\rho(x^2 + y^2)}{lz}} e^{\frac{j2\rho(xx' + yy')}{lz}} dx dy
 \end{aligned}$$

if $e^{-\frac{j\rho(x^2 + y^2)}{lz}} \gg 1$ then $A(x', y', z=0)$ is the FT of $A(x, y, z)$

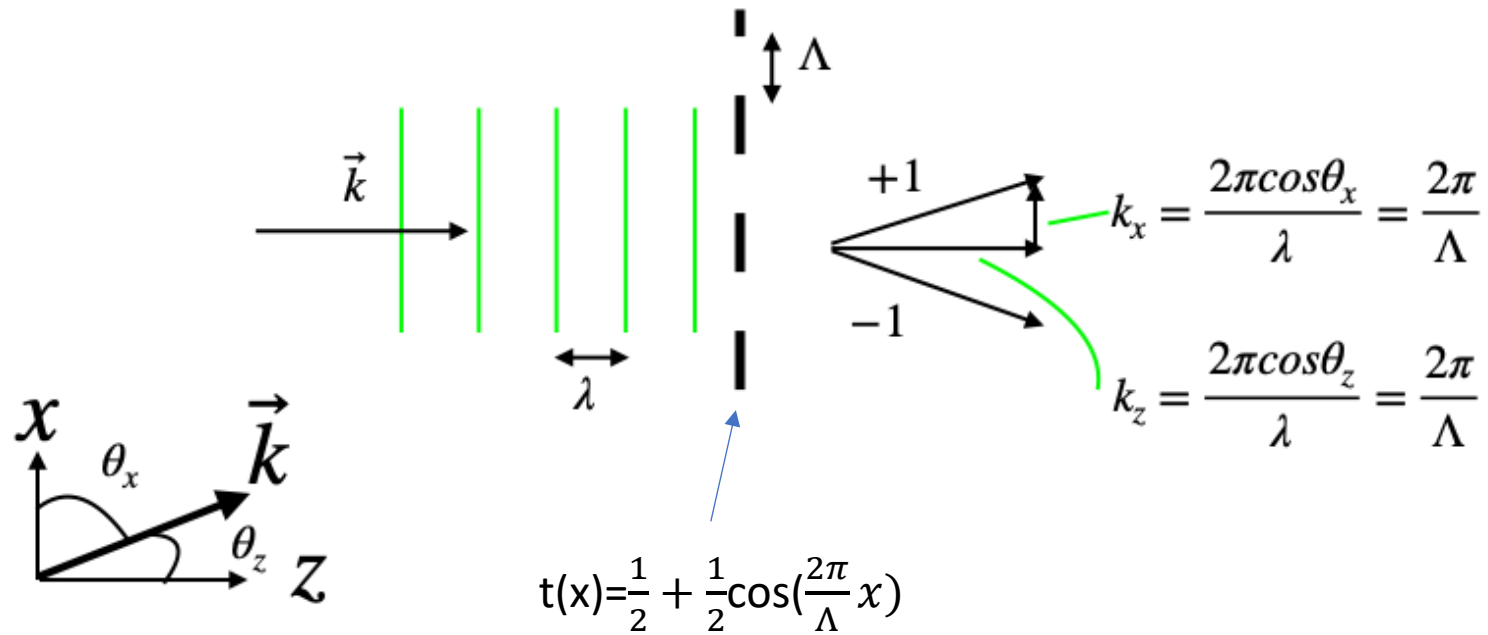
$$u = x / lz \quad v = y / lz$$

Exercise #1

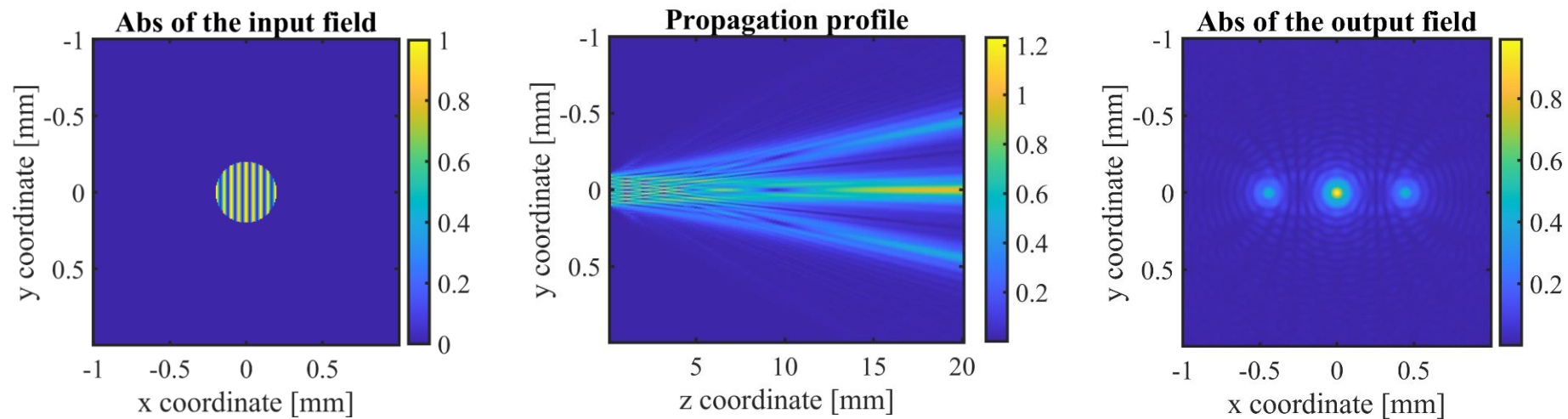
Grating

$$k_x^2 + k_y^2 + k_z^2 = \omega^2 \mu \epsilon = \left(\frac{2\pi}{\lambda}\right)^2 = k^2$$

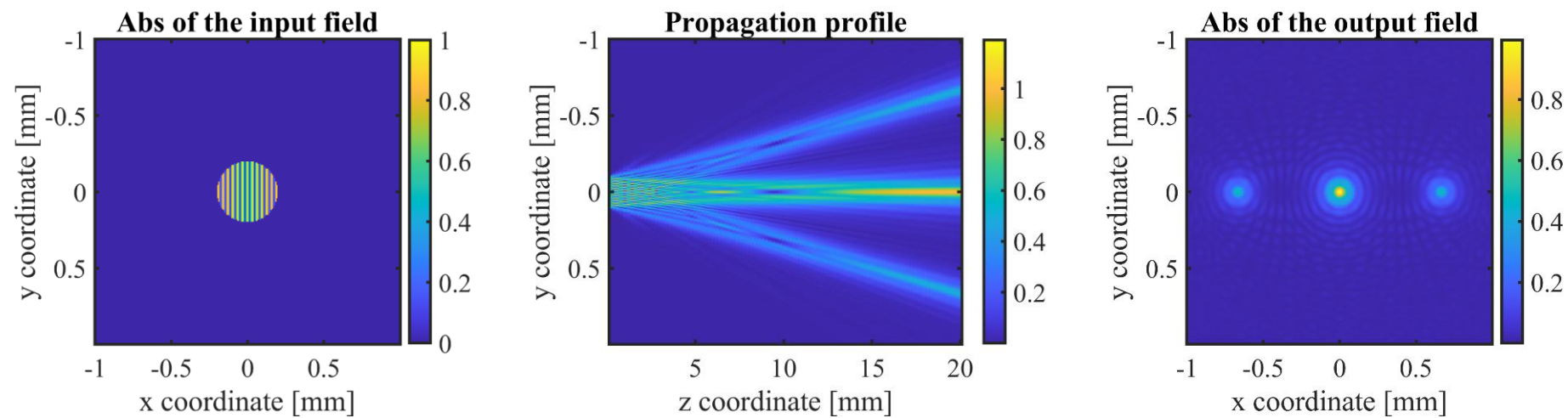
$$k_z = \sqrt{k^2 - (k_x^2 + k_y^2)}$$



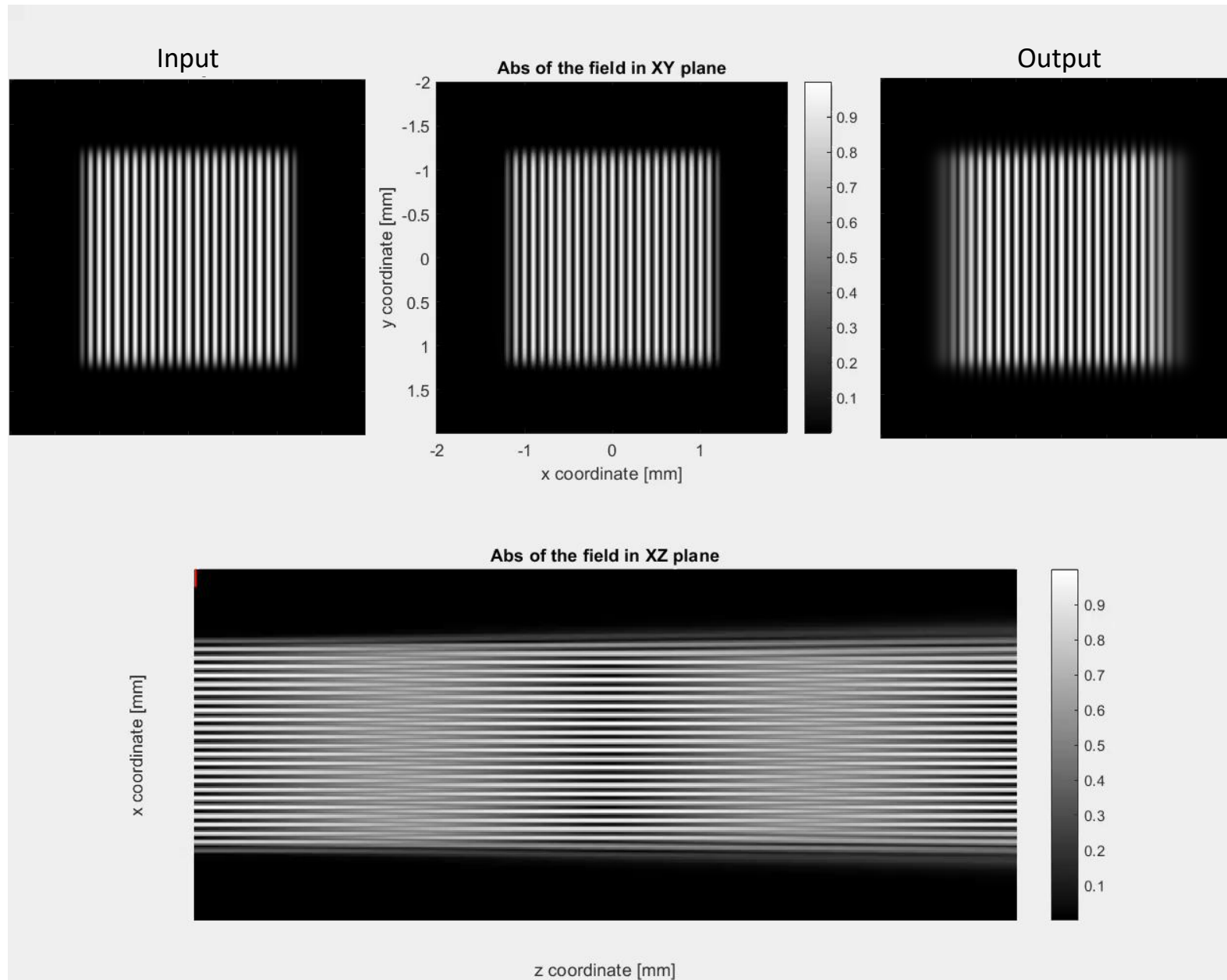
Period: 24 μm



Period: 16 μm



Talbot Effect



Talbot effect explanation

Angular Spectrum

$$E(x, y, z, t) = e^{j\omega t} \iint F(u, v) e^{-j(2pux + 2pvy + \sqrt{k^2 - (2pu)^2 - (2pv)^2} z)} du dv$$

for diffraction order n : $\frac{k_x^2}{2k} z = \frac{(2pn/a)^2}{2(2p/l)} z$

for diffraction order $n+1$: $\frac{k_x^2}{2k} z = \frac{(2p(n+1)/a)^2}{2(2p/l)} z$

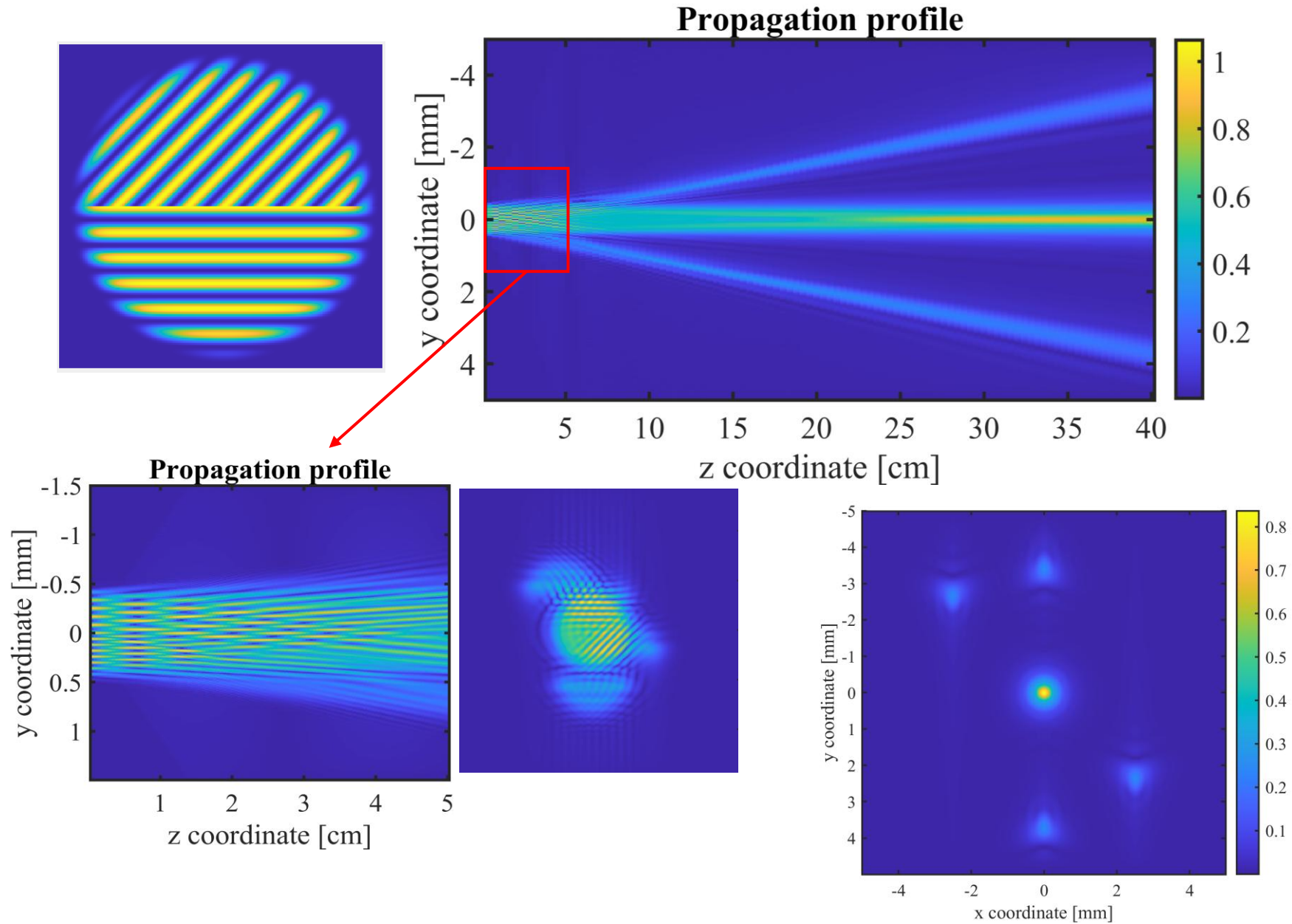
difference : $\frac{(\frac{2p}{a})^2 [(n+1)^2 - n^2] z}{4p/l} = p \frac{l}{a^2} (2n+1) z$

$p z_T = \frac{2a^2}{l}$

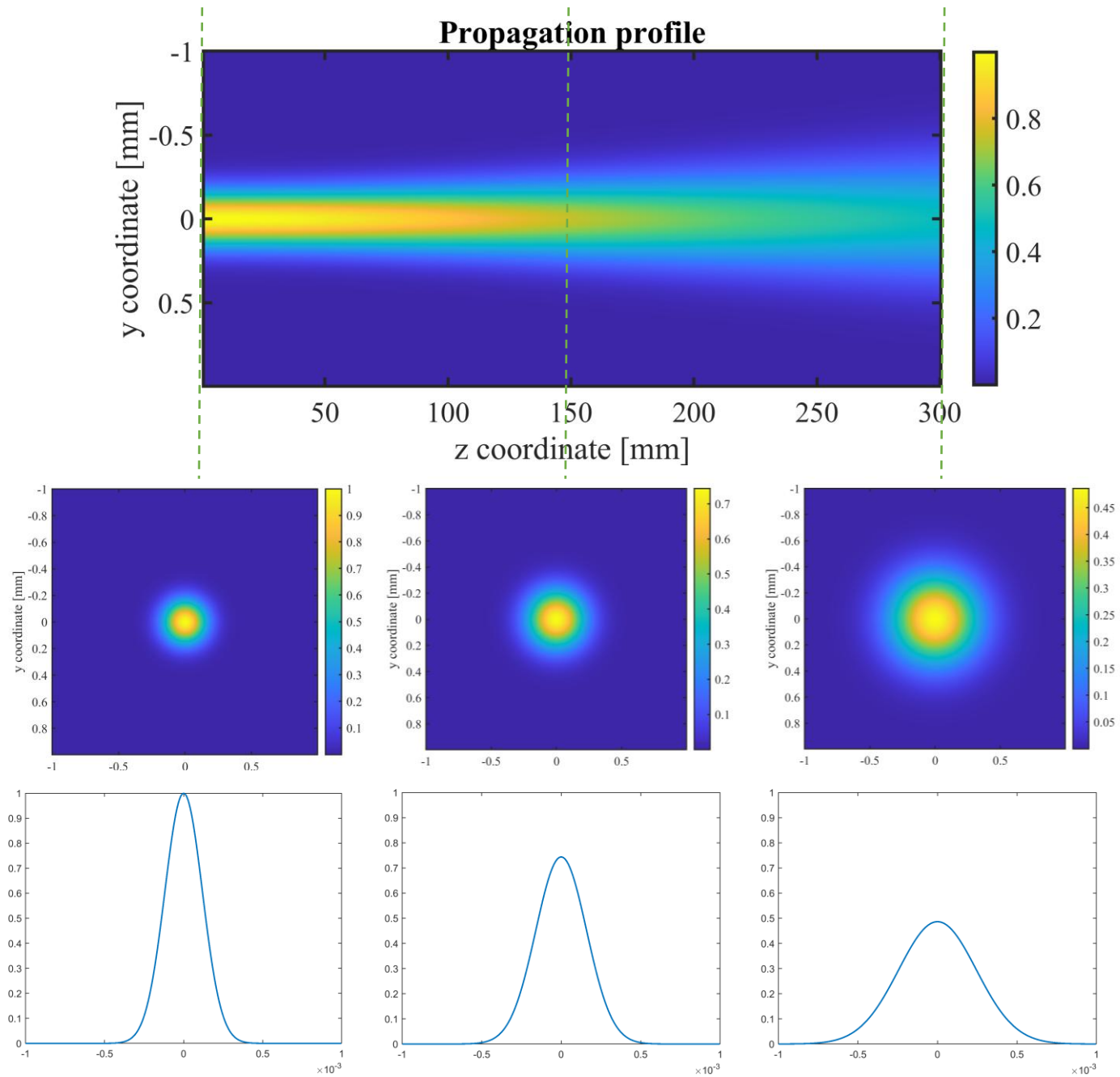




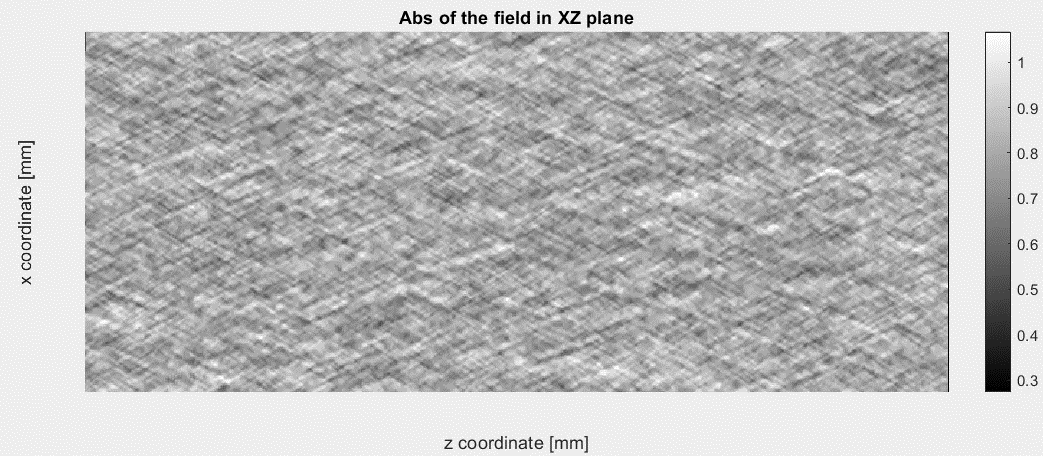
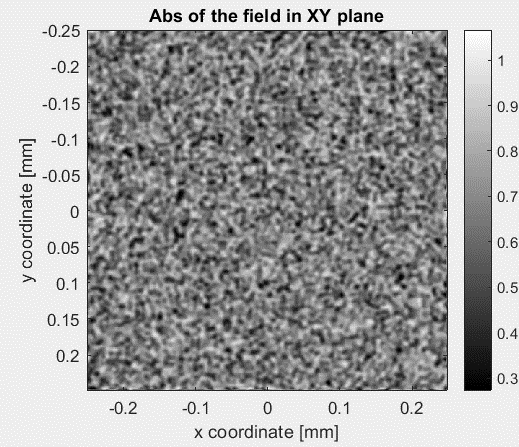
Approaching far field



Gaussian beam



Non-Magic Speckle



Magic speckle

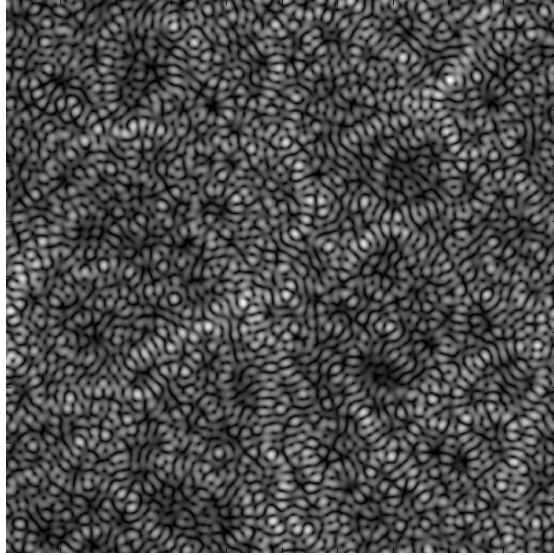
From the paraxial plane wave equation

$$\square A(2pu, 2pv, z) = e^{\frac{-j[(2pu)^2 + (2pv)^2]z}{2k}} \square A(2pu, 2pv, z=0)$$

$$\square A(2pu, 2pv, z) \gg \square A(2pu, 2pv, z=0) \text{ if } u^2/z \ll 1$$

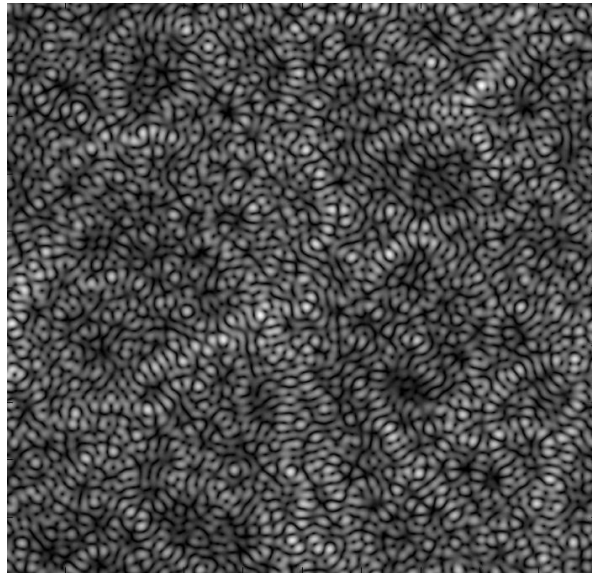
if $(2pu)^2 + (2pv)^2 = \text{const}$

Magic Speckle input

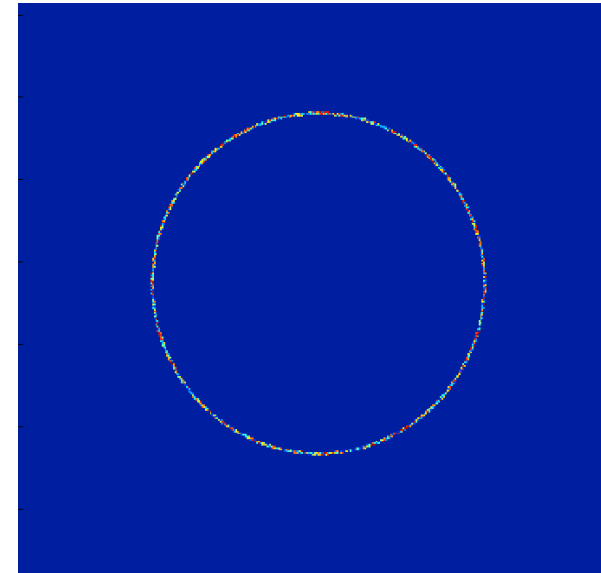


Why doesn't
It change??

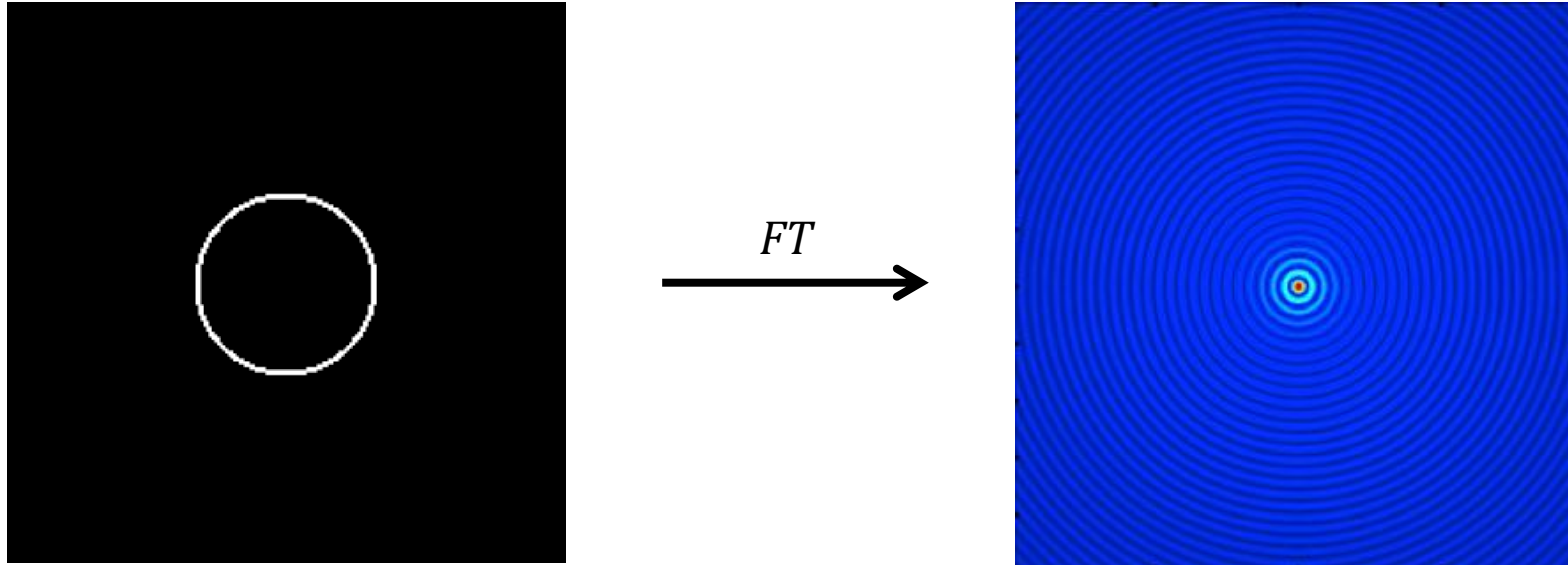
Magic Speckle input



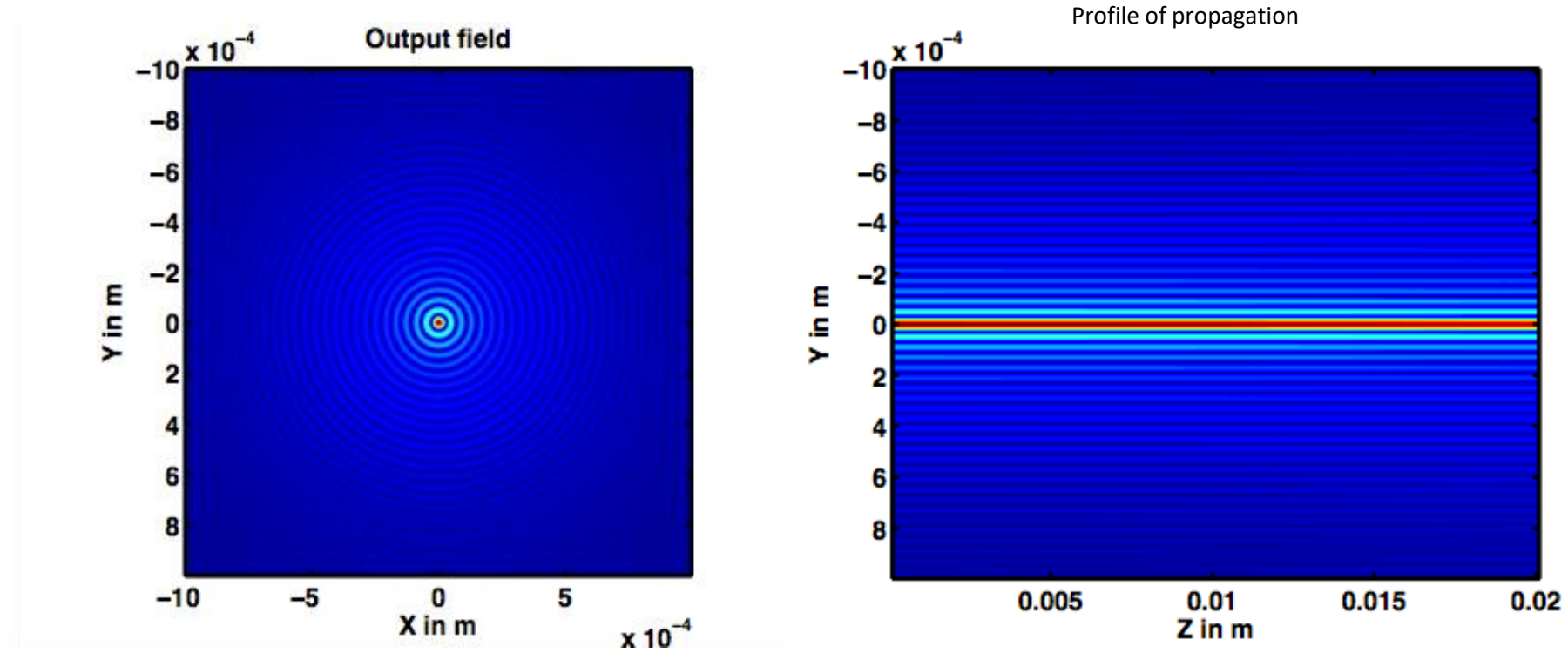
FT



Bessel beam

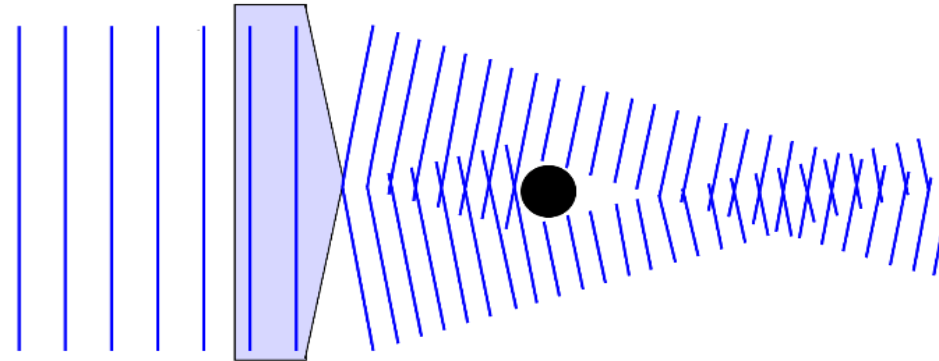
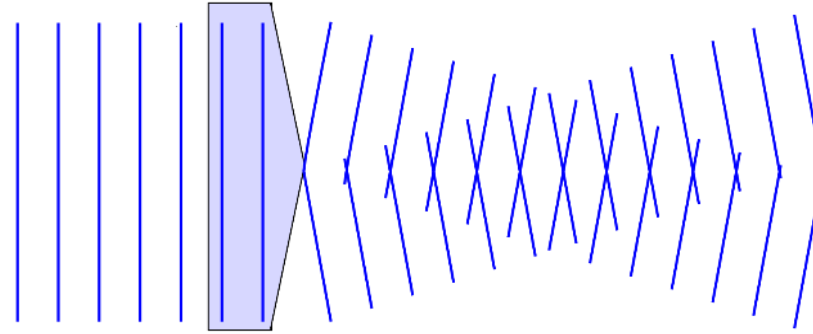


Propagation of a Bessel beam

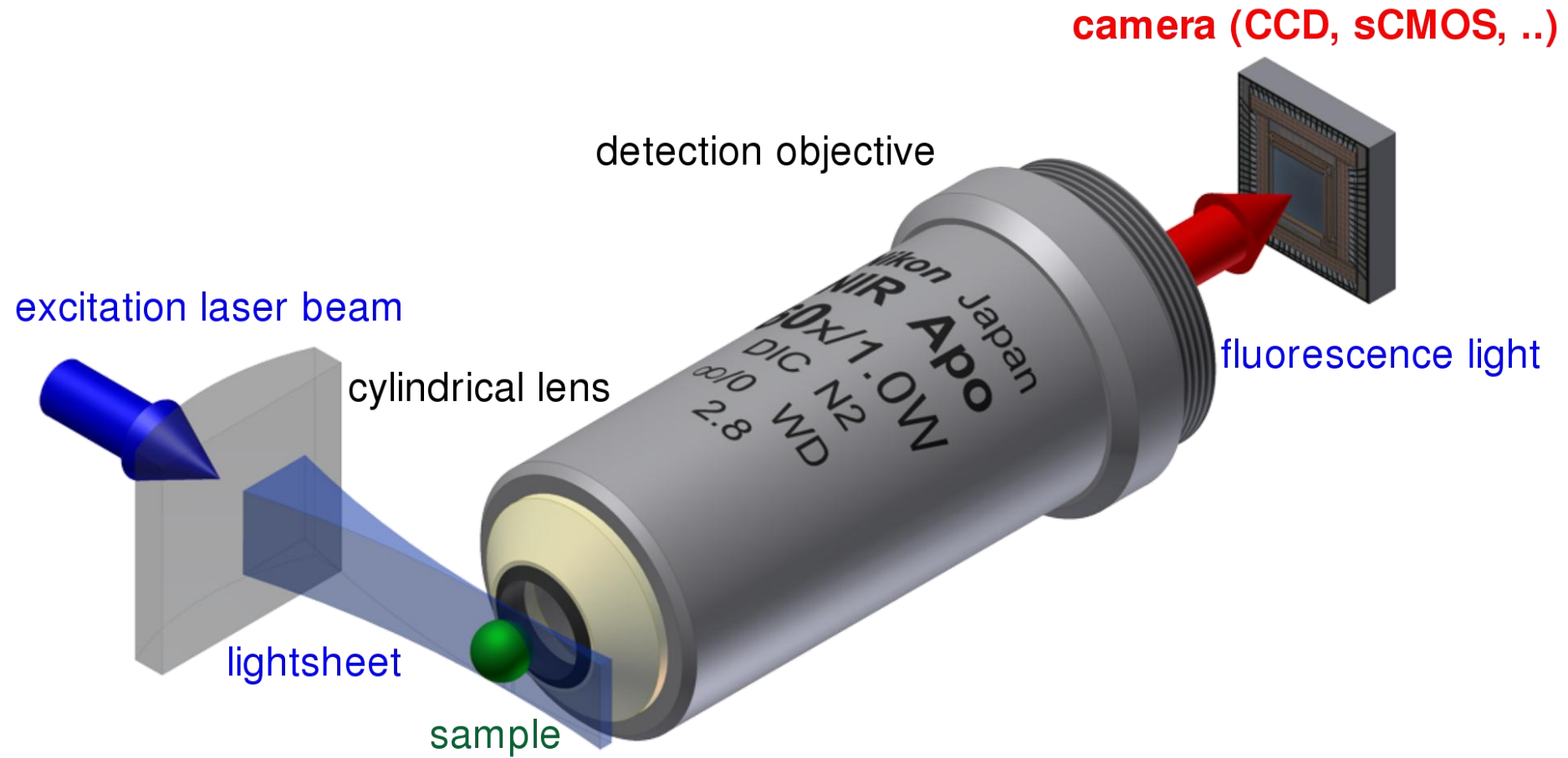


Almost no diffraction is observed after 2cm of propagation in free space

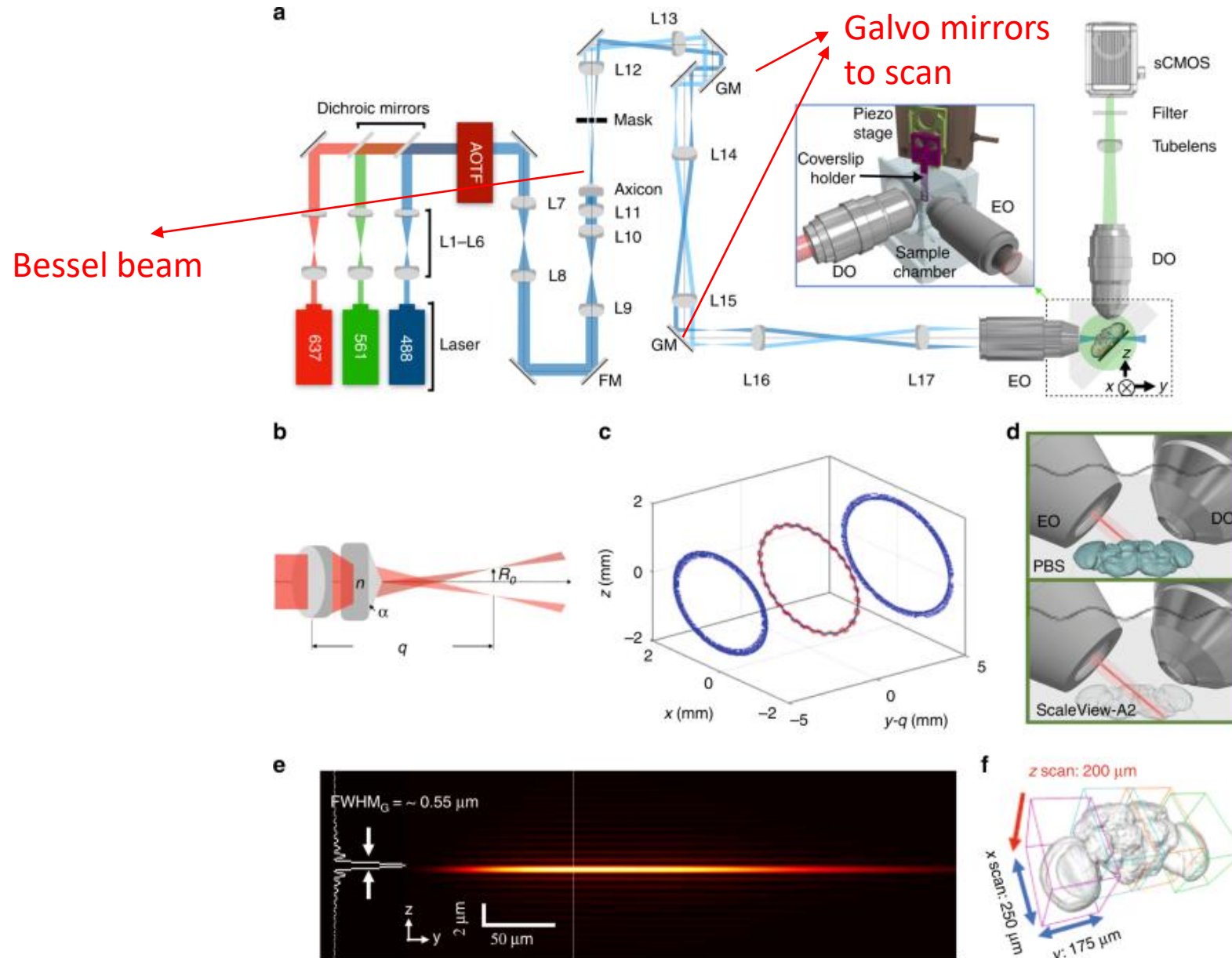
Self healing property



Light Sheet Imaging



Light Sheet by scanning a Bessel beam



Light Sheet by scanning a Bessel beam

